

# Two dimensional polar coordinate system in airy stress functions

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## ABSTRACT

Satisfy the given equations, boundary conditions and biharmonic equation. In order to solve 2-dimensional airy stress function problems by using a polar coordinate reference frame, the equations of equilibrium, definition of Airy's stress function and one of the stress equations of compatibility must be recognized in terms of polar coordinates.

**Keywords:** Form Cartesian Coordinates in to Polar Coordinates, Airy stress function for polar coordinates, Equation in polar coordinates, A stress field symmetric about an axis, A circular hole sheet under remote shear and Example

## 1. INTRODUCTION

The main purpose of this address is to bring to the attention of the workers in Airy stress function and related branches of applied mathematics a simple general method of solution of several important classes of 2-dimensional boundary value problem. A 2-dimensional polar coordinate study in rings and disks, curved bars of narrow rectangular cross section with a circular axis etc. Using the polar coordinates is advantageous to solve in airy stress function.

### 1.1. From Cartesian Coordinates in to Polar Coordinates

To transform equations from Cartesian to polar coordinates, first note the relations

$$x = r \cos \theta,$$

$$y = r \sin \theta$$

$$r^2 = x^2 + y^2,$$

$$\theta = \tan^{-1} \frac{y}{x}$$

$$r = \sqrt{x^2 + y^2},$$

$$\theta = \tan^{-1} \frac{y}{x}$$

$$\frac{\partial r}{\partial x} = \cos \theta$$

$$\frac{\partial \theta}{\partial x} = \frac{-y}{x^2 + y^2} = -\frac{\sin \theta}{r}$$

$$\frac{\partial r}{\partial y} = \sin \theta,$$

$$\frac{\partial \theta}{\partial y} = \frac{x}{x^2 + y^2} = \frac{\cos \theta}{r}$$

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$$\begin{aligned}\frac{\partial \phi}{\partial x} &= \frac{\partial \phi}{\partial r} \frac{\partial r}{\partial x} + \frac{\partial \phi}{\partial \theta} \frac{\partial \theta}{\partial x} \\ \frac{\partial \phi}{\partial x} &= \cos \theta \frac{\partial \phi}{\partial r} - \frac{\sin \theta}{r} \frac{\partial \phi}{\partial \theta} \\ \frac{\partial \phi}{\partial y} &= \frac{\partial \phi}{\partial r} \frac{\partial r}{\partial y} + \frac{\partial \phi}{\partial \theta} \frac{\partial \theta}{\partial y} \\ \frac{\partial \phi}{\partial y} &= \sin \theta \frac{\partial \phi}{\partial r} + \frac{\cos \theta}{r} \frac{\partial \phi}{\partial \theta} \\ \frac{\partial^2 \phi}{\partial x^2} &= \frac{\partial}{\partial x} \left( \frac{\partial \phi}{\partial x} \right) = \\ & \left( \cos \theta \frac{\partial}{\partial r} - \frac{\sin \theta}{r} \frac{\partial}{\partial \theta} \right) \left( \cos \theta \frac{\partial \phi}{\partial r} - \frac{\sin \theta}{r} \frac{\partial \phi}{\partial \theta} \right) \\ \frac{\partial^2 \phi}{\partial x^2} &= (\cos \theta)^2 \frac{\partial^2 \phi}{\partial r^2} + (\sin \theta)^2 \left( \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} \right) + \\ & \sin 2\theta \left( \frac{1}{r^2} \frac{\partial \phi}{\partial \theta} - \frac{1}{r} \frac{\partial^2 \phi}{\partial r \partial \theta} \right)\end{aligned}\tag{1.1.1}$$

Similarly

$$\begin{aligned}\frac{\partial^2 \phi}{\partial y^2} &= (\sin \theta)^2 \frac{\partial^2 \phi}{\partial r^2} + (\cos \theta)^2 \left( \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} \right) - \\ & \sin 2\theta \left( \frac{1}{r^2} \frac{\partial \phi}{\partial \theta} - \frac{1}{r} \frac{\partial^2 \phi}{\partial r \partial \theta} \right)\end{aligned}\tag{1.1.2}$$

$$\begin{aligned}\frac{\partial^2 \phi}{\partial x \partial y} &= -\sin \theta \cos \theta \left( \frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} \right) - \\ & \cos 2\theta \left( \frac{1}{r^2} \frac{\partial \phi}{\partial \theta} - \frac{1}{r} \frac{\partial^2 \phi}{\partial r \partial \theta} \right)\end{aligned}\tag{1.1.3}$$

From (1.1.1) and (1.1.2)

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = \frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2}\tag{A}$$

The Laplace equation is

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = \frac{\partial^2 \phi}{\partial r^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} = 0\tag{B}$$

## 1.2. Airy stress function for polar coordinates

The polar coordinate in plane elasticity problems such as the stresses in circular rings and disks, curved bars of narrow rectangular cross-section with a circular axis, etc. For a two-dimensional polar coordinate system, the solution to plane stress problems involves the determination of in plane stresses. The stress transformation from the Cartesian coordinate to the polar coordinate is

$$\begin{aligned}\sigma_{rr} &= \sigma_x (\cos \theta)^2 + \sigma_y (\sin \theta)^2 + \\ &\quad 2\tau_{xy} \sin \theta \cos \theta \\ \sigma_{\theta\theta} &= \sigma_x (\sin \theta)^2 + \sigma_y (\cos \theta)^2 - \\ &\quad 2\tau_{xy} \sin \theta \cos \theta \\ \tau_{r\theta} &= -\sigma_x \sin \theta \cos \theta + \sigma_y \sin \theta \cos \theta + \\ &\quad \tau_{xy} \left( (\cos \theta)^2 - (\sin \theta)^2 \right)\end{aligned}\quad (1.2.1)$$

## 1.3. Equation in polar coordinates

The airy stress function is a function of the polar coordinate  $\varphi(r, \theta)$  the stress are expressed in term of the airy stress function

$$\sigma_{rr} = \frac{1}{r^2} \frac{\partial^2 \varphi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \varphi}{\partial r}, \quad \sigma_{\theta\theta} = \frac{\partial^2 \varphi}{\partial r^2}, \quad \tau_{r\theta} = -\frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \varphi}{\partial \theta} \right) \quad (1.3.1)$$

The bi harmonic equation is

$$\left( \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} \right) \left( \frac{\partial^2 \varphi}{\partial r^2} + \frac{1}{r} \frac{\partial \varphi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \varphi}{\partial \theta^2} \right) = 0 \quad (C)$$

## 1.4. A stress field symmetric about an axis

Let the airy stress function be  $\varphi(r)$ . The bi harmonic equation is

$$\left( \frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} \right) \left( \frac{d^2 \varphi}{dr^2} + \frac{1}{r} \frac{d\varphi}{dr} \right) = 0 \quad (1.4.1)$$

Each term in this equation has the same dimension in the independent variable  $r$  such an ODE is known as an equi-dimensional equation solution to an equi-dimensional equation is of the form

$$\varphi = r^m \quad (1.4.2)$$

Let equi-dimensional equation is

$$r^2 (r - 3)^2 = 0 \quad (1.4.3)$$

Into the bi harmonic equation, we obtain that the auxiliary equation is

$$\begin{aligned}m^2(m - 3)^2 &= 0. \\ m^2 = 0 \text{ or } (m - 3)^2 &= 0.\end{aligned}$$

The fourth order algebraic equation has double root of 0 and a double root of 3.

The general solution is

$$\varphi(r) = A \log r + Br^3 \log r + Cr^3 + D \quad (1.4.4)$$

Where A, B, C and D are constant of integration.

The components of the stress field are

$$\begin{aligned} \sigma_{rr} &= \frac{1}{r^2} \frac{\partial^2 \varphi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \varphi}{\partial r} = \frac{A}{r^2} + Br(1 + 3 \log r) + 3Cr \\ \sigma_{\theta\theta} &= \frac{\partial^2 \varphi}{\partial r^2} = -\frac{A}{r^2} + Br(5 + 2 \log r) + 6Cr \\ \tau_{r\theta} &= -\frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \varphi}{\partial \theta} \right) = 0 \end{aligned} \quad (1.4.5)$$

The stress field is linear in A, B and C. The contribution due to A and C are familiar they are the same as the lame problem.

For example,

A hole of radius a in an infinite sheet subject to a remote biaxial stress S, the stress field in the sheet is

$$\sigma_{rr} = S \left[ r - \left( \frac{a}{r} \right)^2 \right], \quad \sigma_{\theta\theta} = S \left[ 2r + \left( \frac{a}{r} \right)^2 \right] \quad (D)$$

The stress concentration factor of this hole is 2 we may compare this problem with that of a spherical cavity in an infinite elasticity solid under remote tension

$$\sigma_{rr} = S \left[ r^2 - \left( \frac{a}{r} \right)^3 \right], \quad \sigma_{\theta\theta} = S \left[ 2r^2 + \left( \frac{a}{r} \right)^3 \right] \quad (E)$$

### 1.5. A circular hole in an sheet under remote shear

The sheet is a state of pure shear

$$\tau_{xy} = s, \quad \sigma_{xx} = 0, \quad \sigma_{yy} = 0. \quad (1.5.1)$$

The remote stresses in the polar coordinates are

$$\tau_{r\theta} = S \cos 2\theta, \quad \sigma_{rr} = S \sin 2\theta, \quad \sigma_{\theta\theta} = -S \sin 2\theta \quad (1.5.2)$$

We know that

$$\sigma_{rr} = \frac{1}{r^2} \frac{\partial^2 \varphi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \varphi}{\partial r}, \quad \sigma_{\theta\theta} = \frac{\partial^2 \varphi}{\partial r^2}, \quad \tau_{r\theta} = -\frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \varphi}{\partial \theta} \right)$$

We guess that the stress function must be in the form

$$\varphi(r, \theta) = f(r) \sin 2\theta \quad (1.5.3)$$

The bi harmonic equation is

$$\left( \frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{4}{r^2} \right) \left( \frac{d^2 f}{dr^2} + \frac{1}{r} \frac{df}{dr} - \frac{4f}{r^2} \right) = 0 \quad (F)$$

A solution to this equi-dimensional ODE takes the form

$$f(r) = r^m \quad (1.5.4)$$

Inserting this form into the ODE,

We obtain that

$$\begin{aligned} ((m-2)^2 - 4)((m^2 - 2)) &= 0 \\ ((m-2)^2 - 4) &= 0 \text{ or } (m^2 - 2) = 0 \end{aligned}$$

The algebraic equation has four roots 2, -2, 0, -4.

The stress function is

$$\varphi(r, \theta) = \left( Ar^2 + Br^4 + \frac{C}{r^2} + D \right) \sin 2\theta \quad (1.5.5)$$

The stress components inside the sheet are

$$\begin{aligned} \sigma_{rr} &= \frac{1}{r^2} \frac{\partial^2 \varphi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \varphi}{\partial r} = - \left( 2A + \frac{6C}{r^4} + \frac{4D}{r^2} \right) \sin 2\theta \\ \sigma_{\theta\theta} &= \frac{\partial^2 \varphi}{\partial r^2} = \left( 2A + 12Br^2 + \frac{6C}{r^4} \right) \sin 2\theta \\ \tau_{r\theta} &= - \frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \varphi}{\partial \theta} \right) = \left( -2A - 6Br^2 + \frac{6C}{r^4} + \frac{2D}{r^2} \right) \cos 2\theta \end{aligned} \quad (1.5.6)$$

Where A, B, C and D are constant.

To find the constant A, B, C, D. We invoke the boundary conditions:

1. Remote from the hole namely,  $r \rightarrow \infty$ ,  $\sigma_{rr} = S \sin 2\theta$ ,  $\tau_{r\theta} = S \cos 2\theta$ , giving  $A = \frac{-S}{2}$ ,  $B = 0$ .
2. On the surface of the hole, namely,  $r = a$ ,  $\sigma_{rr} = 0$ ,  $\tau_{r\theta} = 0$ , giving  $D = \frac{-Sa^4}{2}$ .

The stress field inside the sheet is

$$\begin{aligned} \sigma_{rr} &= S \left( 1 + 3 \left[ \frac{a}{r} \right]^4 - 4 \left[ \frac{a}{r} \right]^2 \right) \sin 2\theta \\ \sigma_{\theta\theta} &= -s \left( 1 + 3 \left[ \frac{a}{r} \right]^4 \right) \sin 2\theta \\ \tau_{r\theta} &= s \left( 1 - 3 \left[ \frac{a}{r} \right]^4 + 2 \left[ \frac{a}{r} \right]^2 \right) \cos 2\theta \end{aligned} \quad (1.5.7)$$

**Example 2.1.1** A thin plate is subjected to uniform tensile stress  $\sigma_0$  at its ends; Find the field of stress existing within the plate.

**Solution:**

The origin of coordinate axes at the centre of the plate

The state of stress in the plate is

$$\sigma_x = \sigma_0, \sigma_y = \tau_{xy} = 0$$

The stress function is  $\phi = \sigma_0 \frac{y^2}{2}$ , satisfies the biharmonic equation.

The stress function  $\phi$  may be transformed by substituting  $y = r \sin \theta$ .

$$\phi = \frac{1}{2} \sigma_0 r^2 (\sin \theta)^2$$

$$\phi = \frac{1}{4} \sigma_0 r^2 (1 - \cos 2\theta)$$

The stresses in the plate

$$\sigma_{rr} = \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \phi}{\partial r} = \frac{1}{2} \sigma_0 (1 + \cos 2\theta)$$

$$\sigma_{\theta\theta} = \frac{\partial^2 \phi}{\partial r^2} = \frac{1}{2} \sigma_0 (1 - \cos 2\theta)$$

$$\tau_{r\theta} = -\frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) = -\frac{1}{2} \sigma_0 \sin 2\theta$$

**Example 2.1.2** A thin plate containing a small circular hole of radius  $\alpha$  is subjected to simple tension. Find the field of stress and compare with those of a plate containing a small circular hole.

**Solution:** The boundary conditions appropriate to the circumference of the hole are

$$\sigma_r = \tau_{r\theta} = 0, r = a \quad (2.1.1)$$

For large distance away from the origin we set  $\sigma_r$ ,  $\sigma_\theta$  and  $\tau_{r\theta}$  equal to the values found for a solid plate for  $r = \infty$ .

$$\sigma_r = \frac{1}{2} \sigma_0 (1 + \cos 2\theta)$$

$$\sigma_\theta = \frac{1}{2} \sigma_0 (1 - \cos 2\theta)$$

$$\tau_{r\theta} = -\frac{1}{2} \sigma_0 \sin 2\theta \quad (2.1.2)$$

We assume a stress function

$$\phi = f_1(r) + f_2(r) \cos 2\theta \quad (2.1.3)$$

In  $f_1$  and  $f_2$  are yet to be determined

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr}\right) \left(\frac{d^2 f_1}{dr^2} + \frac{1}{r} \frac{df_1}{dr}\right) = 0 \quad (2.1.4)$$

$$\left(\frac{d^2}{dr^2} + \frac{1}{r} \frac{d}{dr} - \frac{4}{r^2}\right) \left(\frac{d^2 f_2}{dr^2} + \frac{1}{r} \frac{df_2}{dr} - \frac{4f_2}{r^2}\right) = 0 \quad (2.1.5)$$

The solution of equation is

$$f_1 = Ar^2 \log r + B \log r + Cr^2 + D \quad (2.1.6)$$

$$f_2 = Er^2 + Fr^4 + \frac{G}{r^2} + H \quad (2.1.7)$$

Where A, B, C, D, E, F, G and H are constant of integration. The stress function is then obtain integrating equation (2.1.6) and (2.1.7) into (2.1.3) by substituting  $\varphi$  into

The stress are found to be

$$\begin{aligned} \sigma_{rr} &= \frac{1}{r^2} \frac{\partial^2 \varphi}{\partial \theta^2} + \frac{1}{r} \frac{\partial \varphi}{\partial r} = A(1 + 2 \log r) + 2B + C \frac{1}{r} - \left(2E + \frac{6G}{r^4} + \frac{4H}{r^2}\right) \cos 2\theta \\ \sigma_{\theta\theta} &= \frac{\partial^2 \varphi}{\partial r^2} = A(3 + 2 \log r) + 2B - \frac{c}{r^2} + \left(2E + 6Fr^2 + \frac{6G}{r^4}\right) \cos 2\theta \\ \tau_{r\theta} &= -\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \varphi}{\partial \theta}\right) = \left(2E + 6Fr^2 - \frac{6G}{r^4} - \frac{2H}{r^2}\right) \sin 2\theta \end{aligned} \quad (2.1.8)$$

The absence of D indicates that it has no influence on the solution.

According to the boundary condition (2)  $A = F = 0$  in equation (8) because as  $r \rightarrow \infty$  the stresses must assume finite values. Then, according to the condition (a), the equation (2.1.8) yield

$$2B + \frac{C}{a^2} = 0, \quad 2E + \frac{6G}{a^4} + \frac{4H}{a^2} = 0, \quad 2E - \frac{6G}{a^4} + \frac{2H}{a^2} = 0 \quad (G)$$

Also, from equation (2.1.2) and (2.1.8)

We have

$$\sigma_r = -4F, \quad \sigma_0 = 4B \quad (H)$$

Solving the preceding five expressions, we obtain

$$B = \frac{\sigma_0}{4}, \quad C = \frac{-a^2 \sigma_0}{2}, \quad E = \frac{-\sigma_0}{4}, \quad G = \frac{-a^4 \sigma_0}{4}, \quad H = \frac{a^2 \sigma_0}{2} \quad (I)$$

The determination of the stress distribution in a large plate containing a small circular hole is completed by substituting these constant into equation (2.1.8)

$$\sigma_r = \frac{1}{2} \sigma_0 \left[ \left(1 - \left(\frac{a}{r}\right)^2\right) + \left(1 + 3\left(\frac{a}{r}\right)^4 - 4\left(\frac{a}{r}\right)^2\right) \cos 2\theta \right]$$

$$\sigma_{\theta} = \frac{1}{2} \sigma_0 \left[ \left( 1 + \left( \frac{a}{r} \right)^2 \right) - \left( 1 + 3 \left( \frac{a}{r} \right)^4 \right) \cos 2\theta \right]$$

$$\tau_{r\theta} = -\frac{1}{2} \sigma_0 \left[ 1 - 3 \left( \frac{a}{r} \right)^4 + 2 \left( \frac{a}{r} \right)^2 \right] \sin 2\theta \quad (2.1.9)$$

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