# ON CONSTRUCTING SOME MEMBRANES FOR A SYMMETRIC $\alpha$-STABLE PROCESS 

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#### Abstract

Two kinds of membranes located on a fixed hyperplane $S$ in a Euclidean space are constructed for a symmetric $\alpha$-stable process with $\alpha \in(1,2)$. The first one has the property of killing the process at the points of the hyperplane with some given intensity $(r(x))_{x \in S}$. This kind of membranes can be called an elastic screen for the process, by analogy to that in the theory of diffusion processes. The second one has the property of delaying the process at the points of $S$ with some given coefficient $(p(x))_{x \in S}$. In other words, the points of $S$, where $p(x)>0$, are sticky for the process constructed. We show that each one of the membranes is associated with some initial-boundary value problem for pseudo-differential equations related to a symmetric $\alpha$-stable process.


## 1. Introduction

Let $\left(x(t), \mathcal{M}_{t}, \mathbb{P}_{x}\right)$ be a standard Markov process in a $d$-dimensional Euclidean space $\mathbb{R}^{d}$ whose transition probability density $g_{0}$ (with respect to the Lebesgue measure on $\mathbb{R}^{d}$ ) is given by the equality

$$
\begin{equation*}
g_{0}(t, x, y)=(2 \pi)^{-d} \int_{\mathbb{R}^{d}} \exp \left\{i(x-y, \xi)-c t|\xi|^{\alpha}\right\} d \xi, \quad t>0, x \in \mathbb{R}^{d}, y \in \mathbb{R}^{d} \tag{1.1}
\end{equation*}
$$

where $c>0$ and $\alpha \in(1,2)$ are fixed parameters (see [4, Theorem 3.14]). This process is called a symmetric (more precisely, rotationally invariant) $\alpha$-stable process. The generator of it is denoted by $\mathbf{A}$ and this is a pseudo-differential operator with its symbol given by $\left(-c|\xi|^{\alpha}\right)_{\xi \in \mathbb{R}^{d}}$.

Let $\nu$ be a fixed unit vector in $\mathbb{R}^{d}$ and $S$ denote the hyperplane in $\mathbb{R}^{d}$ orthogonal to $\nu$, that is $S=\left\{x \in \mathbb{R}^{d}:(x, \nu)=0\right\}$. By $\mathbf{B}_{\nu}$ we denote a pseudo-differential operator with the function $\left(2 i c|\xi|^{\alpha-2}(\xi, \nu)\right)_{\xi \in \mathbb{R}^{d}}$ as its symbol.

We will consider two kinds of transformations of the process $(x(t))_{t \geq 0}$ (this is a short notation for our process). The first one is connected with the Feynman-Kac formula. Let $(r(x))_{x \in S}$ be a given bounded continuous function with non-negative values. We show that there exists a W-functional $\left(\eta_{t}(r)\right)_{t \geq 0}$ of the process $(x(t))_{t \geq 0}$

[^0]such that its characteristic is given by
\[

$$
\begin{equation*}
\mathbb{E}_{x} \eta_{t}(r)=\int_{0}^{t} d \tau \int_{S} g_{0}(\tau, x, y) r(y) d \sigma_{y}, \quad t \geq 0, x \in \mathbb{R}^{d} \tag{1.2}
\end{equation*}
$$

\]

where the inner integral is a surface one.
As is well-known (see [4, Chapter 10]), there exists a standard Markov process $\left(x^{*}(t), \mathcal{M}_{t}^{*}, \mathbb{P}_{x}^{*}, \zeta\right)$ in $\mathbb{R}^{d}(\zeta$ is the life time of the process) such that the equality

$$
\begin{equation*}
\mathbb{E}_{x}^{*}\left(\varphi\left(x^{*}(t)\right) \mathbb{I}_{\zeta>t}\right)=\mathbb{E}_{x}\left(\varphi(x(t)) \exp \left\{-\eta_{t}(r)\right\}\right) \tag{1.3}
\end{equation*}
$$

is valid for $t>0, x \in \mathbb{R}^{d}$ and $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$ (this is the notation for the Banach space of all continuous bounded functions on $\mathbb{R}^{d}$ with real values and the norm $\left.\|\varphi\|=\sup _{x \in \mathbb{R}^{d}}|\varphi(x)|\right)$. We show that the function (1.3) (denote it by $u(t, x, \varphi)$, $\left.t \geq 0, x \in \mathbb{R}^{d}, \varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)\right)$ is a solution to the following initial-boundary value problem.

Problem $A$. For a given $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$, a continuous function $(u(t, x))_{t>0, x \in \mathbb{R}^{d}}$ is being looked for such that it satisfies
(i) the equation $\frac{\partial u}{\partial t}=\mathbf{A} u$ in the region $t>0, x \notin S$;
(ii) the initial condition $u(0+, x)=\varphi(x)$ for all $x \in \mathbb{R}^{d}$;
(iii) the boundary condition $\frac{1}{2} \mathbf{B}_{\nu} u(t, \cdot)(x+)-\frac{1}{2} \mathbf{B}_{\nu} u(t, \cdot)(x-)=r(x) u(t, x)$ for all $t>0$ and $x \in S$.
The symbol $\mathbf{B}_{\nu} u(t, \cdot)(x+)$ (respectively, $\left.\mathbf{B}_{\nu} u(t, \cdot)(x-)\right)$ for $t>0$ and $x \in S$ means the limit value of the function $\mathbf{B}_{\nu} u(t, \cdot)(z)$, as $z$ approaches $x$ along any curve lying in a finite closed cone $\mathcal{K}$ in $\mathbb{R}^{d}$ with vertex at $x$ such that $\mathcal{K} \subset\{z \in$ $\left.\mathbb{R}^{d}:(z, \nu)>0\right\} \cup\{x\}$ (respectively, $\left.\mathcal{K} \subset\left\{z \in \mathbb{R}^{d}:(z, \nu)<0\right\} \cup\{x\}\right)$.

The second transformation is connected with some random change of time. Let a continuous bounded function $(p(x))_{x \in S}$ with non-negative values be given. For $t \geq 0$, we put

$$
\zeta_{t}=\inf \left\{s \geq 0: s+\eta_{s}(p) \geq t\right\}, \quad \hat{x}(t)=x\left(\zeta_{t}\right), \quad \hat{\mathcal{M}}_{t}=\mathcal{M}_{\zeta_{t}}
$$

It is well-known (see, for example, $[4$, Chapter 10$]$ ) that the process $\left(\hat{x}(t), \hat{\mathcal{M}}_{t}, \mathbb{P}_{x}\right)$ is also a standard Markov process in $\mathbb{R}^{d}$. We show that the function

$$
\begin{equation*}
\hat{u}(t, x, \varphi)=\mathbb{E}_{x} \varphi(\hat{x}(t)), \quad t \geq 0, x \in \mathbb{R}^{d} \tag{1.4}
\end{equation*}
$$

is a solution to the following problem.
Problem B. For a given $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$, a continuous function $(u(t, x))_{t>0, x \in \mathbb{R}^{d}}$ is being looked for such that it satisfies the condition (i), the initial condition (ii) and the following boundary condition (for $t>0$ and $x \in S$ )
(iii') $p(x) \frac{\partial u}{\partial t}(t, x)=\frac{1}{2} \mathbf{B}_{\nu} u(t, \cdot)(x+)-\frac{1}{2} \mathbf{B}_{\nu} u(t, \cdot)(x-)$.
If $\alpha=2$ (and $c=\frac{1}{2}$ ), then our process is a standard Brownian motion, and the operator $\mathbf{A}$ coincides with $\frac{1}{2} \Delta$ ( $\Delta$ is the Laplace operator) and $\mathbf{B}_{\nu}$ coincides with $\frac{\partial}{\partial \nu}$ (the derivative in the direction $\nu$ ). The facts that in this case the functions (1.3) and (1.4) solve Problems $A$ and $B$, respectively, are well-known (some results of the kind can be found in the books $[4,6]$ and also in $[1,2,7]$ and many others).

The article is organized as follows. In Section 2 some auxiliary results are presented. Sections 3 and 4 are devoted to solving the Problems $A$ and $B$, respectively.

## 2. Single-layer Potentials for a Symmetric $\alpha$-stable Process and the Feynman-Kac Formula.

2.1. The function $g_{0}$ defined by (1.1) is continuous in the region $t>0, x \in \mathbb{R}^{d}$ and $y \in \mathbb{R}^{d}$. Moreover, it is uniformly continuous in any region of the form $(t, x, y) \in[\gamma, \infty) \times \mathbb{R}^{d} \times \mathbb{R}^{d}$ for $\gamma>0$. As follows from [3], it satisfies the inequality

$$
\begin{equation*}
g_{0}(t, x, y) \leq N \frac{t}{\left(t^{1 / \alpha}+|y-x|\right)^{d+\alpha}}, \quad t>0, x \in \mathbb{R}^{d}, y \in \mathbb{R}^{d} \tag{2.1}
\end{equation*}
$$

where $N$ is a positive constant. The inequalities of the kind in more general situations including similar inequalities for (fractional) derivatives of $g_{0}$ can be found in [5].
2.2. Let $\nu \in \mathbb{R}^{d}$ be a fixed unit vector and $S$ be the hyperplane in $\mathbb{R}^{d}$ orthogonal to $\nu$. The following formula

$$
\begin{equation*}
\int_{S} e^{i(\xi, y)} g_{0}(t, x, y) d \sigma_{y}=\frac{1}{\pi} \int_{0}^{\infty} e^{-c t\left(|\xi|^{2}+\rho^{2}\right)^{\alpha / 2}} \cos (\rho(x, \nu)) d \rho \tag{2.2}
\end{equation*}
$$

holds true for all $t>0, x \in \mathbb{R}^{d}$ and $\xi \in S$ (see [8]). Combining (2.1) and (2.2) (for $\xi=0$ ), we arrive at the inequality

$$
\begin{equation*}
\int_{S} g_{0}(t, x, y) d \sigma_{y} \leq N \frac{t}{\left(t^{1 / \alpha}+|(x, \nu)|\right)^{1+\alpha}} \tag{2.3}
\end{equation*}
$$

valid for all $t>0$ and $x \in \mathbb{R}^{d}$ with some positive constant $N$.
2.3. In accordance with the definition of $\mathbf{B}_{\nu}$ (see Section 1 ), the following equality (for fixed $t>0$ and $y \in \mathbb{R}^{d}$ )

$$
\mathbf{B}_{\nu} g_{0}(t, \cdot, y)(x)=\frac{2 i c}{(2 \pi)^{d}} \int_{\mathbb{R}^{d}} \exp \left\{i(x-y, \xi)-c t|\xi|^{\alpha}\right\}|\xi|^{\alpha-2}(\xi, \nu) d \xi
$$

is fulfilled for all $x \in \mathbb{R}^{d}$. Integrating by parts leads us to the formula

$$
\begin{equation*}
\mathbf{B}_{\nu} g_{0}(t, \cdot, y)(x)=\frac{2(y-x, \nu)}{\alpha t} g_{0}(t, x, y) \tag{2.4}
\end{equation*}
$$

2.4. Let $(\psi(t, x))_{t \geq 0, x \in S}$ be a continuous function with real values satisfying the inequality $|\psi(t, x)| \leq C t^{-\beta}$ for all $t>0$ and $x \in S$ with some constants $C>0$ and $\beta<1$. We put

$$
\begin{equation*}
V_{0}(t, x)=\int_{0}^{t} d \tau \int_{S} g_{0}(t-\tau, x, y) \psi(\tau, y) d \sigma_{y}, \quad t>0, x \in \mathbb{R}^{d} \tag{2.5}
\end{equation*}
$$

This function is well-defined, as the following estimations show

$$
\begin{aligned}
\left|V_{0}(t, x)\right| & \leq C \int_{0}^{t} \tau^{-\beta} d \tau \int_{S} g_{0}(t-\tau, x, y) d \sigma_{y} \\
& \leq C N \int_{0}^{t} \tau^{-\beta}(t-\tau)^{-1 / \alpha} d \tau \\
& =C N \frac{\Gamma(1-\beta) \Gamma(1-1 / \alpha)}{\Gamma(2-\beta-1 / \alpha)} t^{1-\beta-1 / \alpha}
\end{aligned}
$$

Moreover, this function is continuous in the region $t>0$ and $x \in \mathbb{R}^{d}$. It is called a single-layer potential.

The following properties of the function $V_{0}$ are proved in [8].
2.4.A. The function $V_{0}$ is a solution of the equation $\frac{\partial V_{0}}{\partial t}=\mathbf{A} V_{0}$ in the region $t>0$ and $x \notin S$.
2.4.B. The following relations $\mathbf{B}_{\nu} V_{0}(t, \cdot)(x \pm)=\mp \psi(t, x)$ are held for all $t>0$ and $x \in S$ (the sense of the left hand side is explained in Section 1).

Remark 2.1. Relation 2.4.B are some analogy to the well-known theorem on the jump of the (co-)normal derivative of a single-layer potential in the classical theory of potentials. The term analogous to the so-called direct value of the derivative vanishes in 2.4.B, since $\mathbf{B}_{\nu} g_{0}(t, \cdot, y)(x)=0$ for $y \in S$ and $x \in S$ (see (2.4)).
2.5. Let $(v(x))_{x \in \mathbb{R}^{d}}$ be a continuous bounded function with real values. We put for $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right), t>0$ and $x \in \mathbb{R}^{d}$

$$
Q(t, x, \varphi)=\mathbb{E}_{x}\left(\varphi(x(t)) \exp \left\{\int_{0}^{t} v(x(\tau)) d \tau\right\}\right)
$$

The well-known Feynman-Kac formula asserts that $Q$ satisfies the equation

$$
\frac{\partial Q}{\partial t}=\mathbf{A} Q+v(x) Q
$$

in the region $(t, x) \in(0, \infty) \times \mathbb{R}^{d}$ and the initial condition $Q(0+, x, \varphi)=\varphi(x)$ for all $x \in \mathbb{R}^{d}$.

An intermediate stage of this result is the following integral equation for $Q$

$$
Q(t, x, \varphi)=\int_{\mathbb{R}^{d}} g_{0}(t, x, y) \varphi(y) d y+\int_{0}^{t} d \tau \int_{\mathbb{R}^{d}} g_{0}(t-\tau, x, y) Q(\tau, y, \varphi) v(y) d y
$$

where $t>0, x \in \mathbb{R}^{d}$.

## 3. Solving Problem A

3.1. Let the hyperplane $S$ and the bounded continuous function $(r(x))_{x \in S}$ be such as above. One can easily verify that the function

$$
f_{t}(x)=\int_{0}^{t} d \tau \int_{S} g_{0}(t-\tau, x, y) r(y) d \sigma_{y}
$$

is a W -function for the process $(x(t))_{t \geq 0}$ (see [4, Chapter $\left.6, \S 3\right]$ ) satisfying the inequality

$$
f_{t}(x) \leq N\|r\| \frac{\alpha}{\alpha-1} t^{1-1 / \alpha}
$$

for all $t \geq 0$ and $x \in \mathbb{R}^{d}$ (see (2.4)), where $\|r\|=\sup _{x \in S} r(x)$. Therefore, according to Theorem 6.6 from [4], there exists a W-functional $\left(\eta_{t}(r)\right)_{t \geq 0}$ of the process $(x(t))_{t \geq 0}$ such that $\mathbb{E}_{x} \eta_{t}(r)=f_{t}(x)$ for all $t \geq 0$ and $x \in \mathbb{R}^{d}$.

For $r_{0}(x) \equiv 1$ we put $\eta_{t}=\eta_{t}\left(r_{0}\right), t \geq 0$. The functional $\left(\eta_{t}\right)_{t \geq 0}$ is called the local time on $S$ for the process $(x(t))_{t \geq 0}$. It is evident that $\eta_{t}(r)=\int_{0}^{t} r(x(s)) d \eta_{s}$, $t \geq 0$.
3.2. We now approximate the functional $\left(\eta_{t}(r)\right)_{t \geq 0}$ by somewhat simpler ones. For $h>0$, we define a function $\left(v_{h}(x)\right)_{x \in \mathbb{R}^{d}}$ by setting $v_{h}(x)=\int_{S} g_{0}(h, x, y) r(y) d \sigma_{y}$, $x \in \mathbb{R}^{d}$, and a functional $\left(\eta_{t}^{(h)}(r)\right)_{t \geq 0}$ by the equality $\eta_{t}^{(h)}(r)=\int_{0}^{t} v_{h}(x(s)) d s$, $t \geq 0$.

The function $v_{h}$ for fixed $h>0$ is continuous and bounded, so the W-functional $\left(\eta_{t}^{(h)}(r)\right)_{t \geq 0}$ is well-defined. Its characteristic is given by

$$
\begin{aligned}
f_{t}^{(h)}(x) & =\mathbb{E}_{x} \eta_{t}^{(h)}(r)=\int_{0}^{t} d \tau \int_{\mathbb{R}^{d}} g_{0}(\tau, x, y) v_{h}(y) d y \\
& =\int_{h}^{t+h} d \tau \int_{S} g_{0}(\tau, x, y) r(y) d \sigma_{y}
\end{aligned}
$$

Hence,
$\mathbb{E}_{x} \eta_{t}^{(h)}(r)-\mathbb{E}_{x} \eta_{t}(r)=\int_{t}^{t+h} d \tau \int_{S} g_{0}(\tau, x, y) r(y) d \sigma_{y}-\int_{0}^{h} d \tau \int_{S} g_{0}(\tau, x, y) r(y) d \sigma_{y}$.
Taking into account (2.4), we arrive at the inequality

$$
\begin{aligned}
\sup _{0 \leq t \leq T} \sup _{x \in \mathbb{R}^{d}} \mid \mathbb{E}_{x} \eta_{t}^{(h)}(r) & -\mathbb{E}_{x} \eta_{t}(r) \left\lvert\, \leq N\|r\| \frac{\alpha}{\alpha-1}\left[h^{1-1 / \alpha}\right.\right. \\
& \left.+\sup _{0 \leq t \leq T}\left((t+h)^{1-1 / \alpha}-t^{1-1 / \alpha}\right)\right]
\end{aligned}
$$

valid for all $T>0$ and $h>0$. Denote by $q_{T}(h)$ the expression on the right-hand side of this inequality. Obviously, $q_{T}(h) \rightarrow 0$, as $h \rightarrow 0+$, for any fixed $T>0$. According to Lemma 6.5 from [4], the following inequality

$$
\mathbb{E}_{x}\left(\eta_{t}^{(h)}(r)-\eta_{t}(r)\right)^{2} \leq 2\left(f_{t}^{(h)}(x)+f_{t}(x)\right) q_{T}(h)
$$

holds true for all $t \in[0, T]$ and $x \in \mathbb{R}^{d}$. Since for those $(t, x)$ we have

$$
f_{t}^{(h)}(x) \leq N\|r\| \frac{\alpha}{\alpha-1}(T+h)^{1-1 / \alpha} ; \quad f_{t}(x) \leq N\|r\| \frac{\alpha}{\alpha-1} T^{1-1 / \alpha}
$$

we can assert that the inequality

$$
\begin{equation*}
\mathbb{E}_{x}\left(\eta_{t}^{(h)}(r)-\eta_{t}(r)\right)^{2} \leq 4 N\|r\| \frac{\alpha}{\alpha-1}\left(T+h_{0}\right)^{1-1 / \alpha} q_{T}(h) \tag{3.1}
\end{equation*}
$$

is fulfilled for all $t \in[0, T], x \in \mathbb{R}^{d}$ and $h \in\left(0, h_{0}\right]\left(T>0\right.$ and $h_{0}>0$ are arbitrary fixed numbers).
3.3. For $t>0, x \in \mathbb{R}^{d}$ and $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$, we put

$$
u^{(h)}(t, x, \varphi)=\mathbb{E}_{x}\left(\varphi(x(t)) e^{-\eta_{t}^{(h)}(r)}\right), \quad u(t, x, \varphi)=\mathbb{E}_{x}\left(\varphi(x(t)) e^{-\eta_{t}(r)}\right)
$$

Proposition 3.1. There exists a sequence $\left(h_{n}\right)_{n \geq 1}$ such that $h_{n} \rightarrow 0$, as $n \rightarrow \infty$, and

$$
\lim _{n \rightarrow \infty} u^{\left(h_{n}\right)}(t, x, \varphi)=u(t, x, \varphi)
$$

uniformly with respect to $x \in \mathbb{R}^{d}$ and locally uniformly with respect to $t \in[0, \infty)$.
Proof. Since $\left|e^{-a}-e^{-b}\right| \leq|a-b|$ for all $a \geq 0$ and $b \geq 0$, we can write down the chain of inequalities (for an arbitrary $T>0$ )

$$
\begin{aligned}
& \left|u^{(h)}(t, x, \varphi)-u(t, x, \varphi)\right| \leq\|\varphi\| \mathbb{E}_{x}\left|\eta_{t}^{(h)}(r)-\eta_{t}(r)\right| \\
& \leq\|\varphi\|\left[\mathbb{E}_{x}\left(\eta_{t}^{(h)}(r)-\eta_{t}(r)\right)^{2}\right]^{1 / 2} \leq K_{T}\left(h_{0}\right)\left(q_{T}(h)\right)^{1 / 2}\|\varphi\|
\end{aligned}
$$

valid for all $(t, x) \in[0, T] \times \mathbb{R}^{d}$ and $h \in\left(0, h_{0}\right]$, where $K_{T}\left(h_{0}\right)$ is a constant finite for $T<\infty$. To complete the proof one should make use of the diagonal method.
3.4. The function $u^{(h)}$ (for a fixed $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$ ) is a unique bounded solution to the integral equation (see Section 2.5)
$u^{(h)}(t, x, \varphi)=\int_{\mathbb{R}^{d}} g_{0}(t, x, y) \varphi(y) d y-\int_{0}^{t} d \tau \int_{\mathbb{R}^{d}} g_{0}(t-\tau, x, y) u^{(h)}(\tau, y, \varphi) v_{h}(y) d y$.
It is an easy exercise to verify that the relation

$$
\begin{equation*}
\lim _{h \rightarrow 0+} \int_{\mathbb{R}^{d}} \psi(y) v_{h}(y) d y=\int_{S} \psi(y) r(y) d \sigma_{y} \tag{3.3}
\end{equation*}
$$

is fulfilled for any continuous function $(\psi(y))_{y \in \mathbb{R}^{d}}$ such that $\int_{\mathbb{R}^{d}}|\psi(y)| d y<\infty$.
Proposition 3.2. For a given $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$, the function $(u(t, x, \varphi))_{t \geq 0, x \in \mathbb{R}^{d}}$ is a unique bounded solution of the equation

$$
\begin{equation*}
u(t, x, \varphi)=\int_{\mathbb{R}^{d}} g_{0}(t, x, y) \varphi(y) d y-\int_{0}^{t} d \tau \int_{S} g_{0}(t-\tau, x, y) u(\tau, y, \varphi) r(y) d \sigma_{y} \tag{3.4}
\end{equation*}
$$

Proof. In order to pass to the limit, as $h_{n} \rightarrow 0$, in equation (3.2) (written for $h=h_{n}$ ), one should observe that

$$
\begin{array}{r}
\lim _{n \rightarrow \infty} \int_{0}^{t} d \tau \int_{\mathbb{R}^{d}} g_{0}(t-\tau, x, y) u(\tau, y, \varphi) v_{h_{n}}(y) d y \\
\quad=\int_{0}^{t} d \tau \int_{S} g_{0}(t-\tau, x, y) u(\tau, y, \varphi) r(y) d \sigma_{y}
\end{array}
$$

according to (3.3). Besides,

$$
\int_{0}^{t} d \tau \int_{\mathbb{R}^{d}} g_{0}(t-\tau, x, y) v_{h}(y) d y=f_{t}^{(h)}(x) \leq N\|r\| \frac{\alpha}{\alpha-1}(T+h)^{1-1 / \alpha},
$$

as was established in Section 3.2. Taking into account Proposition 3.1, we arrive at equation (3.4) for the function $u$.

A solution to the equation (3.4) can be constructed by the method of successive approximations. If we put

$$
u_{0}(t, x, \varphi)=\int_{\mathbb{R}^{d}} g_{0}(t, x, y) \varphi(y) d y, \quad t>0, x \in \mathbb{R}^{d}, \varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)
$$

and for $k \geq 1$

$$
u_{k}(t, x, \varphi)=\int_{0}^{t} d \tau \int_{S} g_{0}(t-\tau, x, y) u_{k-1}(\tau, y, \varphi) r(y) d \sigma_{y}
$$

then by induction on $k$, we can easily obtain the following estimate

$$
\begin{equation*}
\left|u_{k}(t, x, \varphi)\right| \leq \frac{\|\varphi\|\|r\|^{k}}{\left(c^{1 / \alpha} \alpha \sin \frac{\pi}{\alpha}\right)^{k}} \frac{t^{k(1-1 / \alpha)}}{\Gamma(k(1-1 / \alpha)+1)} \tag{3.5}
\end{equation*}
$$

held true for all $t>0, x \in \mathbb{R}^{d}, \varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$ and $k=0,1,2, \ldots$ As a consequence of (3.5), we have that the series

$$
\begin{equation*}
\sum_{k=0}^{\infty}(-1)^{k} u_{k}(t, x, \varphi) \tag{3.6}
\end{equation*}
$$

is a continuous solution of (3.4) satisfying the condition

$$
\sup _{t \in[0, T]} \sup _{x \in \mathbb{R}^{d}}|u(t, x, \varphi)|<\infty
$$

for any $T>0$. Another consequence of (3.5) is that such a solution is unique. Therefore, the function $u$ can be represented by the series (3.6). The proposition is proved.
3.5. We now can formulate the main result of Section 3

Theorem 3.3. For a fixed $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$ the function

$$
u(t, x, \varphi)=\mathbb{E}_{x}\left(\varphi(x(t)) e^{-\eta_{t}(r)}\right), \quad t \geq 0, x \in \mathbb{R}^{d}
$$

solves the Problem A.
Proof. The first item on the right hand side of (3.4) satisfies the equation (i) in the whole region $t>0$ and $x \in \mathbb{R}^{d}$. It also satisfies the initial condition (ii). The second item on the right-hand side of (3.4) is a single-layer potential. According to 2.4.A, it satisfies (i) and its initial value vanishes. The relations 2.4.B imply now the equalities

$$
\mathbf{B}_{\nu} u(t, \cdot, \varphi)(x \pm)=\frac{2}{\alpha t} \int_{\mathbb{R}^{d}}(y, \nu) \varphi(y) g_{0}(t, x, y) d y \pm r(x) u(t, x, \varphi)
$$

valid for $t>0$ and $x \in S$, and the condition (iii) follows from these relations immediately. The theorem has been proved.
3.6. If $d=1$, then $S=\{0\}$ and $r=r(0)$ is a non-negative number. The equation for the function $u$ in this case can be written as follows

$$
\begin{equation*}
u(t, x, \varphi)=\int_{\mathbb{R}^{1}} g_{0}(t, x, y) \varphi(y) d y-r \int_{0}^{t} g_{0}(t-\tau, x, 0) u(\tau, 0, \varphi) d \tau \tag{3.7}
\end{equation*}
$$

Denote by $\tilde{u}$ and $\tilde{g}_{0}$ the Laplace transformations of the functions $u$ and $g_{0}$, respectively $(\lambda>0)$

$$
\tilde{u}(\lambda, x, \varphi)=\int_{0}^{\infty} u(t, x, \varphi) e^{-\lambda t} d t, \quad \tilde{g}_{0}(\lambda, x, y)=\int_{0}^{\infty} g_{0}(t, x, y) e^{-\lambda t} d t
$$

Then (3.7) implies the equality

$$
\tilde{u}(\lambda, x, \varphi)=\int_{\mathbb{R}^{1}}\left[\tilde{g}_{0}(\lambda, x, y)-\frac{r \tilde{g}_{0}(\lambda, x, 0) \tilde{g}_{0}(\lambda, 0, y)}{1+r \tilde{g}_{0}(\lambda, 0,0)}\right] \varphi(y) d y
$$

where $\tilde{g}_{0}(\lambda, 0,0)=\left[c^{1 / \alpha} \alpha \sin \frac{\pi}{\alpha}\right]^{-1} \lambda^{1 / \alpha-1}$. It means that the resolvent kernel $\tilde{g}^{*}(\lambda, x, y)$ of the process $\left(x^{*}(t)\right)_{t \geq 0}$ (see Section 1 ) is given by

$$
\tilde{g}^{*}(\lambda, x, y)=\tilde{g}_{0}(\lambda, x, y)-\frac{r \tilde{g}_{0}(\lambda, x, 0) \tilde{g}_{0}(\lambda, 0, y)}{1+r \tilde{g}_{0}(\lambda, 0,0)}
$$

for $\lambda>0, x \in \mathbb{R}^{1}$ and $y \in \mathbb{R}^{1}$. One can obtain from this equality, in particular, the Laplace transform for the distribution function of $\zeta$ (the life time of the process $\left.\left(x^{*}(t)\right)_{t \geq 0}\right)$

$$
\mathbb{E}_{x}^{*} e^{-\lambda \zeta}=\frac{r \tilde{g}_{0}(\lambda, x, 0)}{1+r \tilde{g}_{0}(\lambda, 0,0)}, \quad x \in \mathbb{R}^{1}, \lambda>0
$$

## 4. Solving Problem B

4.1. We are now given by a continuous bounded function $(p(x))_{x \in S}$ with nonnegative values. Consider the Markov process $\left(\hat{x}(t), \hat{\mathcal{M}}_{t}, \mathbb{P}_{x}\right)$ defined in Section 1. The resolvent operator for this process can be calculated in the following way (see [6, Chapter II, §6])

$$
\begin{align*}
& \mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda t} \varphi(\hat{x}(t)) d t=\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda t} \varphi\left(x\left(\zeta_{t}\right)\right) d t \\
& =\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda\left(t+\eta_{t}(p)\right)} \varphi(x(t)) d t+\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda\left(t+\eta_{t}(p)\right)} \varphi(x(t)) d \eta_{t}(p), \tag{4.1}
\end{align*}
$$

where $x \in \mathbb{R}^{d}, \lambda>0, \varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$ (we have taken into account that the equality $\zeta_{t}=t^{\prime}$ implies $\left.t=t^{\prime}+\eta_{t^{\prime}}(p)\right)$.
4.2. If we put

$$
Q_{\lambda}(t, x, \varphi)=\mathbb{E}_{x}\left(\varphi(x(t)) e^{-\lambda \eta_{t}(p)}\right), \quad t>0, \lambda>0, x \in \mathbb{R}^{d}, \varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)
$$

then in accordance with Section 3, we have the following equation for $Q_{\lambda}$

$$
Q_{\lambda}(t, x, \varphi)=\int_{\mathbb{R}^{d}} g_{0}(t, x, y) \varphi(y) d y-\lambda \int_{0}^{t} d \tau \int_{S} g_{0}(t-\tau, x, y) Q_{\lambda}(\tau, y, \varphi) p(y) d \sigma_{y}
$$

Multiplying both sides of this equation by $e^{-\lambda t}$ and integrating with respect to $t$ over $(0, \infty)$, we get the equation

$$
\begin{equation*}
U_{1}(\lambda, x, \varphi)=\int_{\mathbb{R}^{d}} \tilde{g}_{0}(\lambda, x, y) \varphi(y) d y-\lambda \int_{S} \tilde{g}_{0}(\lambda, x, y) U_{1}(\lambda, y, \varphi) p(y) d \sigma_{y} \tag{4.2}
\end{equation*}
$$

where $\tilde{g}_{0}(\lambda, x, y)=\int_{0}^{\infty} g_{0}(t, x, y) e^{-\lambda t} d t$ and

$$
U_{1}(\lambda, x, \varphi)=\int_{0}^{\infty} Q_{\lambda}(t, x, y) e^{-\lambda t} d t=\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda\left(t+\eta_{t}(p)\right)} \varphi(x(t)) d t
$$

4.3. To calculate the second item on the right hand side of (4.1), we observe that

$$
\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda\left(t+\eta_{t}(p)\right)} \varphi(x(t)) d \eta_{t}(p)=\lim _{h \rightarrow 0+} \mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda\left(t+\eta_{t}(p)\right)} \varphi(x(t)) v_{h}(x(t)) d t
$$

where this time $v_{h}(x)=\int_{S} g_{0}(h, x, y) p(y) d \sigma_{y}, h>0, x \in \mathbb{R}^{d}$. According to Section 4.2, we have

$$
\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda\left(t+\eta_{t}(p)\right)} \varphi(x(t)) v_{h}(x(t)) d t=U_{1}\left(\lambda, x, \varphi \cdot v_{h}\right)
$$

It is a very simple conclusion that for $\lambda>0, x \in \mathbb{R}^{d}$ and $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{d}\right)$, the relation $\lim _{h \rightarrow 0+} U_{1}\left(\lambda, x, \varphi \cdot v_{h}\right)=U_{2}(\lambda, x, \varphi)$ fulfilled, where $U_{2}$ is the solution to the equation

$$
\begin{equation*}
U_{2}(\lambda, x, \varphi)=\int_{S} \tilde{g}_{0}(\lambda, x, y) \varphi(y) p(y) d \sigma_{y}-\lambda \int_{S} \tilde{g}_{0}(\lambda, x, y) U_{2}(\lambda, y, \varphi) p(y) d \sigma_{y} \tag{4.3}
\end{equation*}
$$

4.4. As a consequence of 2.4.B, we have the following relations

$$
\mathbf{B}_{\nu}\left(\int_{S} \tilde{g}_{0}(\lambda, \cdot, y) \tilde{\psi}(\lambda, y) d \sigma_{y}\right)(x \pm)=\mp \tilde{\psi}(\lambda, x)
$$

valid for $\lambda>0, x \in S$ and any continuous function $(\psi(t, x))_{t \geq 0, x \in S}$ such as in Section 2.4. These relations imply the following ones $(x \in S, \lambda>0)$

$$
\begin{aligned}
& \mathbf{B}_{\nu} U_{1}(\lambda, \cdot, \varphi)(x \pm)=\int_{\mathbb{R}^{d}} \mathbf{B}_{\nu} \tilde{g}_{0}(\lambda, \cdot, y)(x) \varphi(y) d y \pm \lambda p(x) U_{1}(\lambda, x, \varphi) \\
& \mathbf{B}_{\nu} U_{2}(\lambda, \cdot, \varphi)(x \pm)=\mp p(x) \varphi(x) \pm \lambda p(x) U_{2}(\lambda, x, \varphi)
\end{aligned}
$$

4.5. We put $U(\lambda, x, \varphi)=U_{1}(\lambda, x, \varphi)+U_{2}(\lambda, x, \varphi)$. Then

$$
\mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda t} \varphi(\hat{x}(t)) d t=U(\lambda, x, \varphi)
$$

It follows from the equations (4.2), (4.3) that the function $U$ satisfies the equation

$$
\mathbf{A} U=\lambda U-\varphi(x)
$$

in the region $x \notin S$. Besides, it satisfies the boundary condition $(\lambda>0, x \in S)$

$$
\frac{1}{2} \mathbf{B}_{\nu} U(\lambda, \cdot, \varphi)(x+)-\frac{1}{2} \mathbf{B}_{\nu} U(\lambda, \cdot, \varphi)(x-)=p(x)(\lambda U(\lambda, x, \varphi)-\varphi(x))
$$

We have thus proved the following assertion

Theorem 4.1. The function

$$
\hat{U}(t, x, \varphi)=\mathbb{E}_{x} \varphi(\hat{x}(t)), \quad t>0, x \in \mathbb{R}^{d}
$$

solves the Problem B.
4.6. If $d=1$, then

$$
\begin{aligned}
& \mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda t} \varphi(\hat{x}(t)) d t=\frac{p \tilde{g}_{0}(\lambda, x, 0)}{1+\lambda p \tilde{g}_{0}(\lambda, 0,0)} \varphi(0) \\
& \quad+\int_{\mathbb{R}^{1}}\left[\tilde{g}_{0}(\lambda, x, y)-\frac{\lambda p \tilde{g}_{0}(\lambda, x, 0) \tilde{g}_{0}(\lambda, 0, y)}{1+\lambda p \tilde{g}_{0}(\lambda, 0,0)}\right] \varphi(y) d y
\end{aligned}
$$

for all $\lambda>0, x \in \mathbb{R}^{1}$ and $\varphi \in \mathbb{C}_{b}\left(\mathbb{R}^{1}\right)$, where $p=p(0)$ is a non-negative number.
In the case of $p \rightarrow \infty$ the point $x=0$ becomes an absorbing one. In this case

$$
\begin{aligned}
& \mathbb{E}_{x} \int_{0}^{\infty} e^{-\lambda t} \varphi\left(\hat{x}_{\infty}(t)\right) d t=\frac{\tilde{g}_{0}(\lambda, x, 0)}{\lambda \tilde{g}_{0}(\lambda, 0,0)} \varphi(0) \\
& +\int_{\mathbb{R}^{1}}\left[\tilde{g}_{0}(\lambda, x, y)-\frac{\tilde{g}_{0}(\lambda, x, 0) \tilde{g}_{0}(\lambda, 0, y)}{\tilde{g}_{0}(\lambda, 0,0)}\right] \varphi(y) d y
\end{aligned}
$$

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